Evaluation of runaway thermal reactions of di-*tert*-butyl peroxide employing calorimetric approaches

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Abstract Di-*tert*-butyl peroxide (DTBP) is an organic peroxide (OP) which has widespread use in the various chemical industries. In the past, thermal runaway reactions of OPs have been caused by their general thermal instability or by reactive incompatibility in storage or operation, which can create potential for thermal decomposition reaction. In this study, differential scanning calorimetry was applied to measure the heat of decomposition reactions, which can contribute to understand the reaction characteristics of DTBP. Vent sizing package 2 was also employed to evaluate rates of increase for temperature and pressure in decomposition reactions, and then the thermokinetic parameters of DTBP were estimated. Finally, hazard characteristics of the gassy system containing DTBP, specifically with respect to thermal criticality, were clearly identified.

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List of symbols

Α	Pre-exponential factor of Arrhenius equation
	(\min^{-1})
С	Concentration (g cm^{-3})
C_0	Initial concentration (g cm^{-3})
C_P	Heat capacity (J $g^{-1} K^{-1}$)
$C_{P,b}$	Heat capacity of container (J $g^{-1} K^{-1}$)
$C_{P,s}$	Heat capacity of sample (J $g^{-1} K^{-1}$)
E_{a}	Activation energy $(kJ mol^{-1})$
h	Overall heat transfer coefficient of ambient
	medium (kJ m ^{-2} K ^{-1} min ^{-1})
$hS_{(tp)}$	Value of hS at transitional point (kJ min ⁻¹ K ⁻¹)
k	Reaction rate constant $(M^{1-n} \min^{-1})$
$M_{\rm b}$	Mass of container (g)
$M_{\rm s}$	Mass of sample (g)
Ν	Reaction order (dimensionless)
$P_{\rm p}$	Maximum pressure of reaction (bar)
q	Exothermic heat of reaction $(J g^{-1})$
$q_{ m g}$	Heat generation rate (kJ min ⁻¹)
$q_{ m r}$	Heat removal rate by cooling medium $(kJ min^{-1})$
$-r_{\rm b}$	Chemical reaction rate (mole $L^{-1} min^{-1}$)
R	Universal gas constant (J mol ^{-1} K ^{-1})
S	External surface area of container (m ²)
t	Reaction time (min)
Т	Temperature of reactant (K)
$T_{\rm a}$	Ambient temperature under cooling system (K)
$T_{\rm a(tp)}$	Ambient temperature up to the transitional point
	(K)
$T_{\rm c}$	Critical ignition or extinction temperature (K)
$T_{\rm CE}$	Critical extinction temperature (K)
$T_{\rm CI}$	Critical ignition temperature (K)

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$T_{\rm c(tp)}$	Transitional point of critical ignition and extinction				
	temperatures (K)				
T_p	Maximum temperature of reaction (K)				
T_{o}	Onset temperature of reaction (K)				
$T_{\rm s}$	Temperature at the steady state, which occurs at				
	the intersection point of curves $q_{\rm g}$ and $q_{\rm r}$ (K)				
T_{SE}	Final stable point of extinction temperature (K)				
T_{SI}	Final stable point of ignition temperature (K)				
T_{SL}	Low stable temperature at the steady state (K)				
$T_{\rm SH}$	High stable temperature at the steady state (K)				
$T_{\mathbf{M}}$	Intermediate temperature of intersection point of				
	curves $q_{\rm g}$ and $q_{\rm r}$ (K)				
V	Volume of reactant (L)				
X_A	Fractional conversion (dimensionless)				

Greek letters

ho	Density of reactant (g cm ⁻³)	
$\Delta H_{ m d}$	Heat of decomposition reaction (J g^{-1}))

$\Delta H_{\rm t}$	Transitional released heat of reaction at reaction
	time t (J g ⁻¹)
ΛH_{1}	Total heat of reaction $(I \sigma^{-1})$

 ΔH_{total} lotal heat of reaction (J g⁻¹)

- ΔT_p Temperature difference between T_0 and T_p (K)
- Φ Phi factor, thermal inertia (dimensionless)

Introduction

In the past, several significant explosion incidents in Taiwanese industries have been related to the thermal instability of organic peroxides (OPs). OPs, having the peroxy oxygen bond (-O-O-), can easily be broken down due to their highly active characteristics. OPs also react with external materials, such as may be kept nearby in storage facilities, in various ways resulting in large discharges of energy [1-5]. Di-tert-butyl peroxide (DTBP) is one commonly used industrial OP, widely employed as a radical initiator for polymerization, as a source for alkoxyl radicals, as a hardener, as a linking agent, and as a fuel combustion additive or used in reforming operations. Because of this variety of uses, it is necessary to understand the various potential exothermic reactions of DTBP. Due to its reactivity and instability, the National Fire Protection Association (NFPA) of USA has recognized DTBP as one of its class III OPs [6, 7].

In this study, the initial stage was to acquire fundamental thermokinetic data and the essential safety parameters of DTBP by differential scanning calorimetry (DSC) and vent sizing package 2 (VSP2). The second stage was to employ that data, specifically the onset temperature of reaction (T_o), heat of decomposition reaction (ΔH_d), activation energy (E_a), maximum temperature of reaction (T_p), maximum increase in temperature (($dT \ dt^{-1})_{max}$), and maximum increase in pressure (($dP \ dt^{-1})_{max}$), to evaluate the thermokinetic parameters of the decomposition reactions of DTBP. Measuring the thermal parameters can be employed by using experimental methods involving isothermal reaction temperature varying removal heat, and the reaction temperature and heat change throughout the reaction time [8–11]. The third stage was to evaluate the critical temperature of DTBP through numerical simulation according to Semenov's thermal explosion theory [12–15]. Finally, we tried to clearly identify the criterion of critical temperature incurred by the gassy system containing DTBP during undesirable situations.

Experimental setup

Experimental samples

Ninety-eight mass% of DTBP was purchased from Aldrich Co., and then stored in a refrigerator at 4 °C. Its density was 0.796 g cm⁻³. Furthermore, 25 mass% of DTBP was prepared in toluene solution for use in our VSP2 experiments.

Experimental apparatus and procedures

DSC

The DSC is considered as an useful tool for evaluating thermal hazards and investigating decomposition heat of reactive chemicals. Non-isothermal programmed screening experiments were performed on a Mettler TA8000 system DSC 821^e apparatus. Each sample was put into a high pressure gold-plated crucible (ME-26732), which was then manually brought to a hermetic seal by use of a special tool equipped with the Mettler DSC and could withstand a high pressure, up to 15 MPa. STAR^e software was used to acquire thermal curves and to assess the thermokinetics [16, 17]. The heating rates chosen for the temperature-programmed ramps were 1, 4, 6, and 10 °C min⁻¹ from 30 to 300 °C to give the reaction processes better thermal equilibrium. In each case, a 4–8 mg test sample of DTBP was prepared.

VSP2

The VSP2, a PC-controlled adiabatic calorimeter manufactured by Fauske & Associates, Inc., was used to obtain thermokinetic and thermal hazard data for DTBP, such as temperature and pressure traces versus time. The heatwait-search (H-W-S) method for detecting the self-heating rate was employed. If the self-heating rate was greater than 0.1 °C min⁻¹, the H-W-S and main heater would be immediately terminated to measure the phenomenon of self-exothermicity. The low heat capacity of the cell ensured that any reaction heat released remained within the tested sample. Thermokinetic data and pressure behavior in the same test cell (112 mL) could usually be tested without any difficult extrapolation to the process scale due to the low phi factor (Φ) of about 1.05–1.32. This low Φ allows for bench-scale simulation of worst-case scenarios, such as incorrect dosing, cooling failure, or external fire conditions [18, 19].

Results and discussion

Thermal decomposition analysis of DTBP by DSC

Table 1 summarizes the thermokinetic data from the DSC STAR^e program for runaway assessment. The thermal curve of the DSC tests is also shown in Fig. 1 and the decomposition data is listed in Table 1. According to various heating rates, the decomposition temperature of DTBP began at about 98–109 °C, and the overall released heat ($\triangle H_{\text{total}}$) is ca. 1,100 J g⁻¹.

Adiabatic analysis of DTBP by VSP2

The VSP2 is capable of showing temperature and pressure rise under adiabatic conditions. The VSP2 test of DTBP is

 Table 1
 Thermal decomposition data for DTBP at various heating rates by DSC

Heating rate/ °C min ⁻¹	Sample mass/ mg	T₀/ °C	<i>T_p</i> / ℃	$\Delta H_{\rm d}/$ J g ⁻¹	$E_{\rm a}/{ m kJ~mol^{-1}}$	lnk
1	7.4	98	165	1,192	132.60	29.73
4	5.0	109	181	1,116	134.15	30.24
6	4.3	106	187	1,158	131.70	29.52
10	4.7	121	198	745	138.13	30.84



Fig. 1 Thermal decomposition of DTBP at various heating rates by DSC



Fig. 2 Measured transitional temperature and pressure of the decomposition reaction of 25 mass% DTBP by means of VSP2

shown in Fig. 2, where the maximal temperature (T_p) and pressure (P_{max}) were 302 °C and 59.8 bar, respectively.

Evaluation of thermokinetic parameters

Several important numerical assumptions were employed to derive the thermokinetic parameters of an exothermic reaction from experimental data on the rate of self-heating under adiabatic conditions:

- (1) The reaction mechanism was assumed to be independent of temperature, which allowed the temperature and concentration dependencies to be dealt separately. The total heat generated was considered to be directly from the adiabatic temperature increase, assuming constant heat capacity.
- (2) Heat generated in the reaction vessel was assumed to correspond to changes in concentration, such that the rate of change of concentration and the rate of heat generation are directly proportional to the rate of temperature increase under adiabatic conditions, where the extent of the reaction equals the temperature increase expressed as a fraction of the total adiabatic temperature increase. The temperature dependence of the reaction rate constant was assumed to obey the Arrhenius equation.
- (3) The dependence of reaction rate on concentration was represented by first order reactions with fractional values applied such that complex mechanisms can be represented by simple overall kinetic expressions.

In an adiabatic reaction system, the fractional conversion (X_A) and reaction temperature (T) at any reaction time (t) can be correlated with the assumption of constant specific heat capacity by the ratio value of $\triangle H_t$ to $\triangle H_{\text{total}}$ from initiation to end as follows:

$$X_A = \Delta H_t / \Delta H_{\text{total}} = [mC_p(T - T_o)] / [mC_p(T_p - T_o)]$$

= $(T - T_o) / (T_p - T_o),$ (1)

where $T_{\rm o}$ and T_p are the initial and maximal temperatures for the overall decomposition reaction of DTBP, respectively. $\Delta H_{\rm t}$ and $\Delta H_{\rm total}$ are the heat of reaction at temperature *T* and T_p , respectively.

Reaction concentration can be correlated as follows:

$$C = C_0(1 - X_A) = C_0(T_p - T)/(T_p - T_o)$$

= $C_0(T_p - T)/\Delta T_p,$ (2)

where C_0 is the initial concentration, C is the concentration at any time t, and ΔT_p equals $T_p - T_0$.

An *n*th-order rate equation of decomposition reaction for DTBP can be expressed as:

$$-r_{\rm b} = -{\rm d}C/{\rm d}t = kC^n. \tag{3}$$

Substituting Eq. 2 into Eq. 3 and combining with the Arrhenius equation yields

$$k^* = C_0^{n-1} A \exp(-E_a/RT)$$

= $(dT/dt) / [(T_p - T)/\Delta T_p]^n \Delta T_p.$ (4)

Taking the natural logarithms on both sides of Eq. 4 gives the following:

$$\ln k^* = \ln \left\{ (dT/dt) / \left[(T_p - T) / \Delta T_p \right]^n \Delta T_p \right\}$$

= $\ln (C_0^{n-1} A) - E_a / RT,$ (5)

where $dT dt^{-1}$ is the rate of temperature increase between the onset temperature (T_o) and the maximum temperature (T_p) in an adiabatic reaction system. The Arrhenius kinetic parameters for the overall reaction can thenceforth be determined from Eq. 5.

Heat evolved during the VSP2 experiment was absorbed by the sample and sample container. This would be expected to the adiabatic temperature increase and the rate of temperature increase, and must be taken into account during data analysis. The Φ (Phi) factor, or thermal inertia, is defined as $\Phi = 1 + M_b C_{p,b}/M_s C_{p,s}$, where M_b , M_s , and C_p represent the mass of the container, the mass of the sample, and the heat capacity of the sample. The Φ value was used to adjust the self-heating rate as well as the observed adiabatic temperature increase. Further, the initial data obtained by H-W-S processes were experimentally and mathematically insignificant, and were therefore excluded in the derivation of the kinetic parameters. The data from runaway conditions were also excluded due to the instability of the reaction system.

Substituting the modified experimental data from VSP2 into Eq. 5 and assuming n = 1, we can plot lnk versus $-1,000 T^{-1}$ and derive truly significant linear correlations, as illustrated in Fig. 3. As such, the assumption of n = 1for reaction order is reasonable. The activation energy (E_a) and pre-exponential factor (*A*) of the Arrhenius equation can be obtained from the slope and intersection of the line at the vertical axis. In this way, the measured and evaluated thermokinetic parameters of decomposition reactions for DTBP were acquired.

Stability criteria and critical runaway temperature for the decomposition reaction of DTBP

Semenov model assumes an uniform temperature distribution within the reaction system. This assumption is very close to the case of a homogeneous system in a container [20]. The general formula for heat generation rate due to volume of reactant (V) is expressed as:

$$q_{\rm g} = qV(-r_{\rm b}),\tag{6}$$

where *q* is exothermic heat of reaction (*q*) and chemical reaction rate is $-r_b$. Expressing the reaction rate by Arrhenius' method and substituting it into Eq. 6, the heat generation rate can be expressed as:

$$q_{\rm g} = q VAC^n \exp(-E_{\rm a}/RT). \tag{7}$$

Similarly, the heat removal rate from the reaction container to the ambient cooling medium is expressed as:

$$q_{\rm r} = hS(T - T_{\rm a}),\tag{8}$$

where the overall heat transfer coefficient of the ambient medium is h, and the external surface area of the container is S. T and T_a are the temperatures inside the reaction container and of the ambient medium under cooling system, respectively.

The overall energy balance in the control sample volume is that the rate of heat accumulation is equivalent to the rate of heat generation minus the rate of heat removal:



Fig. 3 lnk^* versus 1,000 T⁻¹ for reactive rate prediction of 25 mass% DTBP at n = 1

$$\rho V C_P(\mathrm{d}T/\mathrm{d}t) = q_\mathrm{g} - q_\mathrm{r}. \tag{9}$$

The first term on the right-hand side, q_g , represents the heat generation rate for an exothermic reaction and the second term, q_r , is the heat removal rate from this reaction system. Inserting Eqs. 7 and 8 into Eq. 9, the overall energy balance can be expressed as:

$$\rho V C_P(\mathrm{d}T/\mathrm{d}t) = q V A C^n \exp(-E_\mathrm{a}/RT) - h S(T-T_\mathrm{a}). \tag{10}$$

As mentioned above, it is assumed that the overall reaction order of a decomposition reaction of DTBP, n, is tantamount to unity. Thus, substituting Eq. 2 into Eq. 10 with n = 1, this process can be rewritten as:

$$\rho VC_P(\mathrm{d}T/\mathrm{d}t) = q VAC_0[(T_p - T)/(T_p - T_o)]$$

$$\exp(-E_\mathrm{a}/RT) - hS(T - T_\mathrm{a}). \tag{11}$$

Under a steady state, i.e., (dT/dt) = 0, Semenov's sufficient and necessary conditions for a "critical situation of the reaction system" are

$$q_{\rm g}|_{T=T_{\rm c}} = q_{\rm r}|_{T=T_{\rm c}} \tag{12}$$

and

$$\left. \frac{\mathrm{d}q_{\mathrm{g}}}{\mathrm{d}T} \right|_{T=T_{\mathrm{c}}} = \left. \frac{\mathrm{d}q_{\mathrm{r}}}{\mathrm{d}T} \right|_{T=T_{\mathrm{c}}}.$$
(13)

Equation 11 also denotes the reaction system under the conditions of a steady state, i.e., (dT/dt) = 0, when the rate of heat generation of the system is equal to the rate of heat removal. If q_g is greater than q_r , then (dT/dt) > 0. In such a situation, heat accumulation can lead to runaway reactions. Equation 11 can also be written as:

$$qVAC_{0}[(T_{p} - T_{c})/(T_{p} - T_{o})]\exp(-E_{a}/RT_{c}) = hS(T_{c} - T_{a}).$$
(14)

Applying the conditions of Eqs. 12 and 13 gives:

$$\frac{qVAC_0 \exp(-E_a/RT_c)[-1 + E_a(T_p - T_c)/RT_c^2]}{(T_p - T_o)} = hS.$$
(15)

After dividing Eq. 14 by 15, and some rearranging, the temperature in the reaction system can be acquired by Eq. 16:

$$\frac{E_{\rm a}}{RT_{\rm c}^2} = \frac{1}{(T_{\rm c} - T_{\rm a})} + \frac{1}{(T_p - T_{\rm c})}.$$
(16)

Critical temperature (T_c) can be solved as:

$$T_{\rm c} = \frac{(T_p + T_{\rm a}) \pm \sqrt{(T_p + T_{\rm a})^2 - 4[1 + R(T_p - T_{\rm a})/E_{\rm a}]T_pT_{\rm a}}}{2[1 + R(T_p - T_{\rm a})/E_{\rm a}]}.$$
(17)

Equation 17 indicates that the critical temperature (T_c) is a function of the activation energy (E_a) , the ambient



Fig. 4 Balance diagram of heat generation rate q_{g} , and heat removal rate q_{r} , for the decomposition reaction of DTBP

temperature under cooling system (T_a) , and the maximum temperature of reaction (T_p) under adiabatic conditions. Figure 4 shows an example of calculating critical temperatures for the decomposition of DTBP. T_a was set at 363.15 K and the values of E_a were listed, respectively. Two critical temperature points of $T_{\rm c}$ can be obtained from the diagram of heat generation rate and theoretical critical heat removal rate versus temperature. One was the critical extinction temperature $T_{CE} = 442.13$ K and the other was the critical ignition temperature $T_{\rm CI} = 373.63$ K, which were calculated from Eq. 17. Substituting these two values of T_c into Eq. 15, the values of hS can be obtained, which were the overall heat transfer coefficient multiplied by the external surface of the container. The calculated values of hS from this equation were equal to 13.38 and 1.086 kJ min⁻¹ K⁻¹, which correspond to $T_{\rm CE}$ and $T_{\rm CI}$, respectively. By substituting the values of hS into Eq. 14 and iterating, one can obtain another set of temperatures, denoted as T_{SE} and T_{SI} . These T_{SE} and T_{SI} values were the intersection points of the curve q_g with curves q_{r1} and q_{r3} instead of tangent points, respectively. These two points, $T_{\rm SE}$ and $T_{\rm SI}$, can also be obtained graphically from the intersection points of q_g and q_r when the values of hS are equal to 13.38 and 1.086 kJ min⁻¹ K⁻¹, individually.

We found that the value of $T_{\rm SI}$ was 457.27 K and the value of $T_{\rm SE}$ was 363.46 K in this case. As hS = 13.38 kJ min⁻¹ K⁻¹ and the temperature in the reaction system was greater than 442.13 K, i.e., $T > T_{\rm CE}$, the heat removal rate $q_{\rm r1}$ was greater than the heat generation rate $q_{\rm g}$. Therefore, the temperature in the reaction system decreases continually and moves toward point $T_{\rm CE}$. Finally, it ceases at this point, when the temperature in the reaction system was in the range of $T_{\rm SE} < T < T_{\rm CE}$. Similarly, the heat removal rate $q_{\rm g}$. As such, the temperature of the reaction system approached and terminated at point $T_{\rm SE}$. When the

reaction system temperature was lower than $T_{\rm SE}$, the heat removal rate $q_{\rm r1}$ became less than the heat generation rate $q_{\rm g}$. As a result, the temperature of the reaction system increases and goes back to point $T_{\rm SE}$ at the end. Points $T_{\rm CE}$ and $T_{\rm SE}$ were the critical extinction temperature and the final stable extinction temperature, respectively. These two points also represent the temperature that never increased and the temperature that never decreased at hS = 13.38 kJ min⁻¹ K⁻¹, respectively.

Equations 15 and 17 gave values of T_{CI} and hS equal to 373.63 K and 1.086 kJ min⁻¹ K⁻¹, respectively. The intersection point of curves q_g and q_{r3} at this calculation was $T_{\rm SI} = 457.27$ K. When the temperature in the reaction system was less than 373.63 K, i.e., $T < T_{CI}$, the heat removal rate q_{r3} , was less than the heat generation rate q_g . From this we can see that the temperature in the reaction system increases and moves toward $T_{\rm CI}$. When the temperature of the reaction system was in the range of $T_{\rm CI} < T < T_{\rm SI}$, the heat removal rate q_{r3} , was again lower than the heat generation rate $q_{\rm g}$, which leads us to see the temperature of the reaction system increasing from point $T_{\rm CI}$ moving toward to point T_{SI} . Once the temperature was greater than T_{SI} , the heat removal rate q_{r3} was higher than the heat generation rate q_{g} , which leads us to the temperature of the reaction system moving back to point T_{SI} . Points T_{CI} and T_{SI} were the critical ignition temperature and the final stable ignition temperature, respectively. These two points also represent the temperature that never decreased and the temperature of no return at $hS = 1.086 \text{ kJ min}^{-1} \text{ K}^{-1}$, respectively. When the value of hS is located in the range of 1.086 < hS <13.38 kJ min⁻¹ K⁻¹, three intersection points can be obtained between curves q_{r2} and q_g . These three points, denoted as T_{SL} , T_M , and T_{SH} , individually represent the steady-state temperatures at the low, intermediate, and high points. The curve q_{r2} is shown as a dashed line.

As in previous analysis, it was difficult to arrive at the intermediate point $T_{\rm M}$ because of the instability of this point during chemical reaction. We assumed that it was appropriate to start this reaction system exactly at the temperature $T_{\rm M}$; in this way, if some tiny perturbation in the operating conditions was to take the reaction system away from its steady state, it would not move back to $T_{\rm M}$ but rather would finish up either at the low or the high temperature in a steady state, either T_{SL} or T_{SH} , respectively. Thus, a slight increase in temperature would produce a net heat generation that could boost the temperature even higher and conversely a minute drop in temperature would induce a net heat removal, which would cause the temperature to fall even further. In this case, the intermediate steady-state temperature $T_{\rm M}$ was unstable. By contrast, the steady-state temperatures $T_{\rm SL}$ and $T_{\rm SH}$ were stable, such that if the situation were perturbed in any way, it would return there naturally.



Fig. 5 Correlation of evaluated temperatures T_{CI} , T_{CE} , T_{SI} , and T_{SE} versus T_{a} and the temperature of transitional point for the decomposition reaction of DTBP

Figure 5 shows the correlation of evaluated temperatures $T_{\rm CI}$, $T_{\rm CE}$, $T_{\rm SI}$, and $T_{\rm SE}$ versus ambient temperature $T_{\rm a}$ using the above calculation technique for the decomposition reactions of DTBP. These temperatures were in the order $T_{\rm SI} > T_{\rm CE} > T_{\rm CI} > T_{\rm SE}$. Both $T_{\rm SI}$ and $T_{\rm CE}$ decreased, and both $T_{\rm CI}$ and $T_{\rm SE}$ rose gradually with increasing $T_{\rm a}$. As soon as the ambient temperature $T_{\rm a}$ increased to the transitional point $T_{\rm a(tp)}$, all four of these temperatures of this transitional point can be deduced from the following equations:

$$T_{a(tp)} = \frac{E_a T_p}{(E_a + 4RT_p)} \tag{18}$$

and

$$T_{\rm c(tp)} = \frac{E_{\rm a}T_p}{(E_{\rm a} + 2RT_p)}.$$
(19)

With respect to the criteria for stability or instability in a reaction system, the heat generation and removal diagrams do not permit us to conclude stability when

$$\frac{\mathrm{d}q_{\mathrm{g}}}{\mathrm{d}T} < \frac{\mathrm{d}q_{\mathrm{r}}}{\mathrm{d}T}.$$
(20)

This equation is necessary, but by itself is not sufficient condition for stability. Even still, it does allow us to validate that the state will be unstable when

$$\frac{\mathrm{d}q_{\mathrm{g}}}{\mathrm{d}T} > \frac{\mathrm{d}q_{\mathrm{r}}}{\mathrm{d}T}.$$
(21)

If T_s is defined as a stable temperature in a reaction system at the state of $q_g = q_r$ and a small temperature perturbation as $\delta T = T - T_s$ in this state, the reaction system is stable under the following conditions:



Fig. 6 Relationship between T_a and hS at critical ignition and extinction temperatures for the decomposition reaction of DTBP

if
$$\delta T > 0$$
, $q_g < q_r$
or
if $\delta T < 0$, $q_g > q_r$. (22)

Equation 22 gives both the sufficient and necessary conditions for a stable reaction system. This criterion for stability has been evaluated by Lu et al. [21, 22].

From the above definitions and energy-mass balance equations, the stable and unstable criteria of DTBP decomposition reactions can be determined. As long as the temperature of the reaction system occurred over the curve for $T_{\rm SI}$ or under the curve for $T_{\rm SE}$, the system was stable: the temperature in the reaction system would return to either of these two curves eventually. When the reaction temperature fell between the curves for $T_{\rm CI}$ and $T_{\rm CE}$, the system was unstable. Furthermore, there were two stable areas, T_{SL} and T_{SH} , located between two narrow areas, where $T_{\rm SE} < T_{\rm SL} < T_{\rm CI}$ and $T_{\rm CE} < T_{\rm SH} < T_{\rm SI}$, respectively. As the value of $T_{\rm a}$ increased, the values for the $T_{\rm SI}$ and $T_{\rm CE}$ temperature curves decreased, but the values for the $T_{\rm CI}$ and $T_{\rm SE}$ temperature curves increased at the same time. Thus, when the value of T_a increases, these four temperature curves move closer and closer. As soon as the value of $T_{\rm a}$ exceeded 407.92 K, the phenomenon of criticality vanished. These temperature curves coincide with the transitional point. Here, the value of transitional point $T_{\rm c(tp)}$ was 431.40 K.

The required values for hS at critical runaway and stable temperatures could be evaluated by using either Eq. 14 or 15 after the value of T_c was determined from Eq. 17. The calculated result of T_a versus hS at the critical runaway and stable temperatures for the decomposition reaction of DTBP s is shown in Fig. 6, where we see a boomerangshaped area. The solid curve of hS signifies the critical extinction temperature, T_{CE} , and the final stable extinction



Fig. 7 Correlation of evaluated temperatures T_{CI} , T_{CE} versus T_a and *hS* for the decomposition reaction of DTBP



Fig. 8 Correlation of evaluated temperatures T_c versus T_s and hS for the decomposition reaction of DTBP

Table 2 Thermokinetic parameters of 25 mass% DTBP by VSP2

Φ	<i>T_p</i> / °C	P _p / bar	$\Delta T_p / \Delta t /$ °C min ⁻¹	$\Delta P_{\rm p} / \Delta t /$ bar min ⁻¹	$E_{\rm a}$ /kJ mol ⁻¹	A/min ⁻¹
1.87	303	59.7	281	130.5	134.58	3.19E15
1.79	298	56.7	272	112.5	132.41	2.97E14
1.83	301	58.3	274	124.2	133.79	3.07E15

temperature, T_{SE} , while the other solid curve with dot points inside denotes the curve of *hS* at critical ignition temperature, T_{CI} , and the final stable ignition temperature, T_{SI} . Between these two curves is a boomerang-shaped area which contains one unstable temperature, $T_{\rm M}$, and two stable temperatures, $T_{\rm SL}$ and $T_{\rm SH}$ (Fig. 7). When the temperatures $T_{\rm CE}$, $T_{\rm SE}$, $T_{\rm CI}$, and $T_{\rm SI}$ were equivalent, the value of hS is expressed as $hS_{\rm (tp)}$. The value of $hS_{\rm (tp)}$ was equal to 32.66 kJ min⁻¹ K⁻¹. The values of $hS_{\rm (tp)}$ can be calculated by a rearrangement Eq. 15, which was obtained by combining Eqs. 18 and 19.

$$hS_{(\text{tp})} = \frac{qVAC_0(1 + 4RT_p/E_a)\exp[-(2 + E_a/RT_p)]}{(T_p - T_o)}.$$
 (23)

Above this boomerang-shaped area in the diagram, the area contains the temperatures $T_{\rm SL}$ and $T_{\rm SH}$, which were distributed both above and below this zone. From this diagram, we see that the required value of hS in the decomposition reaction of DTBP was strongly affected by variation of $T_{\rm a}$. Once the value of $T_{\rm a}$ was gathered, the value of $T_{\rm CI}$, $T_{\rm CE}$, and hS could all be estimated from the reaction thermokinetic parameters and Eqs. 15–17 mentioned before (Fig. 8).

Conclusions

The heat of the decomposition reaction of 98 mass% DTBP was approximately 1,100 J g⁻¹, as measured by DSC. The thermokinetic parameters of the decomposition reaction, listed in Table 2, were determined from the experimental result of VSP2. The unstable temperature area of the decomposition reaction of DTBP was shown to be enclosed by the curves $T_{\rm CI}/T_{\rm SI}$ and $T_{\rm CE}/T_{\rm SE}$ in Fig. 6. With the exception of this unstable area, the reaction system was otherwise stable.

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